

The **XXZ** model in the attractive regime with boundaries

Chaiho's home assignment

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→ sG model with boundary interactions

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both repulsive and attractive regime through the NLIE
approach
 - Fusion-TBA approach (Klümper-Pearce,
Kuniba-Sakai-JS)
well studied for the repulsive regime
- Try the fusion approach in the attractive regime and see
the consistency of the Ahn-Nepomichie result.

The outline of the talk

- The model and the scaling limit
- The fusion-TBA approach in the repulsive regime
- The fusion-TBA approach in the attractive regime
 - the origin of the difficulty
 - crucial property: Factorization by Q functions

The XXZ model with boundaries

- The Hamiltonian

$$\mathcal{H} = -\frac{1}{2} \left(\sum_{j=1}^{N-1} (\sigma_j^x \sigma_{j+1}^x + \sigma_j^y \sigma_{j+1}^y + \Delta \sigma_j^z \sigma_{j+1}^z) + p \sigma_1^z + p' \sigma_N^z \right)$$

- Parameterization

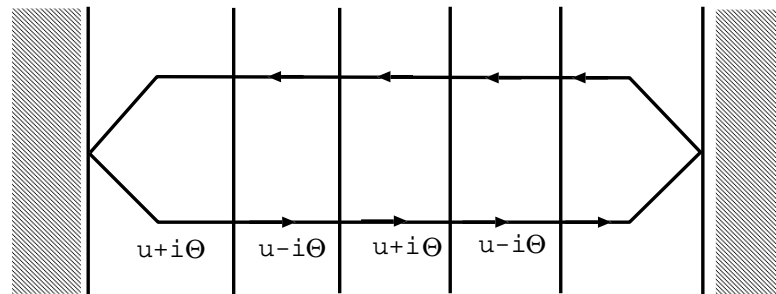
$$\Delta = -\cos \gamma$$

$$p = \sin \gamma \cot \frac{\gamma}{2} (H_+ - 1), \quad p' = \sin \gamma \cot \frac{\gamma}{2} (H_- - 1)$$

- Consider $0 < p < p_C = 1 + \cos \gamma$ (no boundary string) and $p = p'$.

The 6V model with boundaries

- A vertex model with the 6 possible configurations ($N \times N$ sites)
- The Boltzmann weights: parameterized by u, γ, ρ
e.g., $w(+, +, +, +) = \rho \sin(2 + u)\gamma$
- an inhomogeneous model, $u \rightarrow u \pm i\Theta$



The scaling limit

$$N \rightarrow \infty \qquad \Theta \rightarrow \frac{2}{\pi} \log \frac{4N}{mL}$$
$$u \rightarrow \frac{2}{\pi} i\theta$$

Correspondence: $m = \text{sG Kink mass}$, $\beta^2 = 8(\pi - \gamma)$.

The diagonalization of T

- The eigenvalue of the inhomogeneous 6V model

$$T_1(x) = B_-(x)T_0(x+i)\frac{Q(x-2i)}{Q(x)} + B_+(x)T_0(x-i)\frac{Q(x+2i)}{Q(x)}$$

where $u = x/i$ and

$$B_{\pm}(x) = [x \pm iH_+][x \pm iH_-], \quad T_0(x) = [2x][x + \Theta]^N [x - \Theta]^N$$

$$Q(x) = \prod_{j=1}^{N/2} [x - x_j][x + x_j] \quad [x] := \sinh \frac{\gamma}{2}$$

- The energy of the XXZ model

$$E = i \frac{d}{dx} \log T_1(x) \Big|_{x=i}$$

The fusion relations

- Introduce fusion transfer matrices $T_m(x) = \text{Tr}_{V_m} \mathcal{T}(x)$ ($\dim V_m = m + 1$.)
- Functional Relations

$$T_m(x+i)T_m(x-i) = f_m(x) + T_{m+1}(x)T_{m-1}(x)$$

$$f_m(x) = T_0(x + (m+1)i)T_0(x - (m+1)i)$$

$$\times \prod_{j=1}^m B_-(x + (m+2-2j)i)B_+(x - (m+2-2j)i)$$

- The normalization (the ground state):
 - $T_m = \text{degree } N$ in $[x]$ under the p.b.c
 - $T_m = \text{degree } 2N + 2m$ in the case with boundaries

$T - Y$ transformation

- The $T - Y$ transformation (Klümper-Pearce ('91), Kuniba-Nakanishi-JS('94))

$$Y_j(x) = \frac{T_{j-1}(x)T_{j+1}(x)}{f_j(x)} \quad j = 1, 2, \dots$$

- The universal $Y -$ system (Al. B. Zamolodchikov)

$$Y_j(x+i)Y_j(x-i) = (1+Y_{j+1}(x))(1+Y_{j-1}(x)), \quad j = 1, 2, \dots$$

Truncation of relations

- The peculiarity at $\gamma = \pi/(p + 1)$ (BLZ '96)

$$\begin{aligned} T_{p+1}(x) &= T_{p-1}(x)B_-(x - pi)B_+(x + pi) \\ &+ T_0(x + (p + 1)i) \prod_{j=1}^p B_-(x + (p + 1 - 2j)i) \\ &+ T_0(x - (p + 1)i) \prod_{j=1}^p B_+(x + (p + 1 - 2j)i) \end{aligned}$$

or symbolically $T_{p+1}(x) = 2 + T_{p-1}(x)$

- factorization $\gamma = \pi/(p + 1)$ (symbolically)

$$T_p(x + i)T_p(x - i) = (1 + T_{p-1}(x))^2$$

Truncated Y – system

$$Y_j(x+i)Y_j(x-i) = (1 + Y_{j+1}(x))(1 + Y_{j-1}(x)), \quad j = 1, \dots, p-3$$

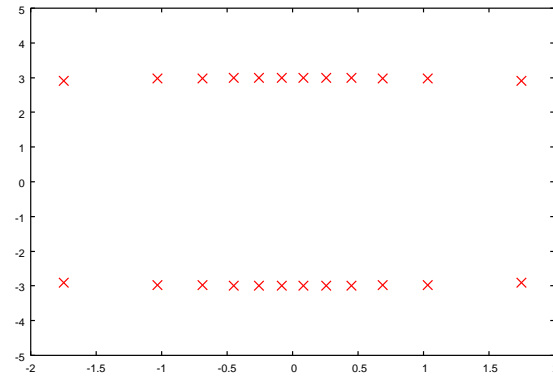
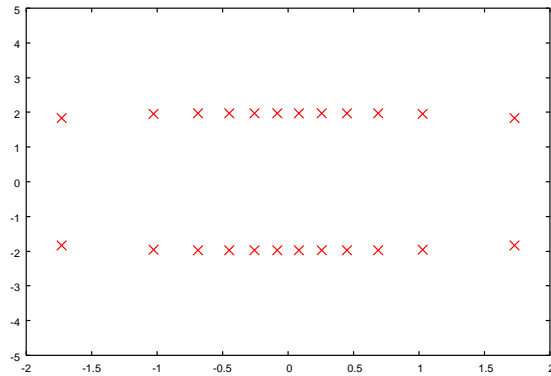
$$Y_{p-2}(x+i)Y_{p-2}(x-i) = (1 + Y_{p-3}(x))(1 + \lambda(x)Y_t(x)) \\ \times (1 + \lambda(\bar{x})Y_t(x))$$

$$Y_{t-1}(x+i)Y_{t-1}(x-i) = (1 + Y_{p-2}(x))$$

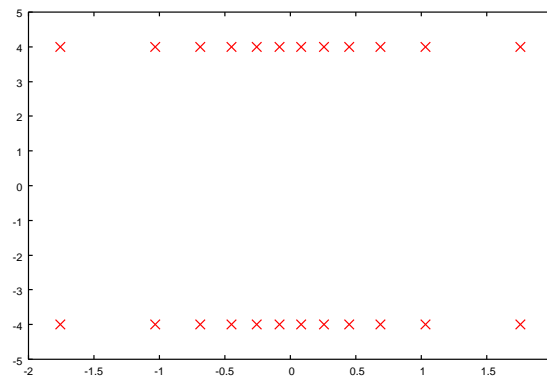
D_p Dynkin type equations.

Zeros of T in the repulsive regime

$$N = 24, p.b.c., \gamma = \frac{\pi}{4}$$



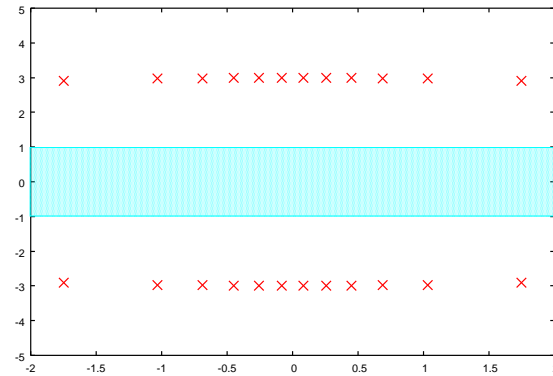
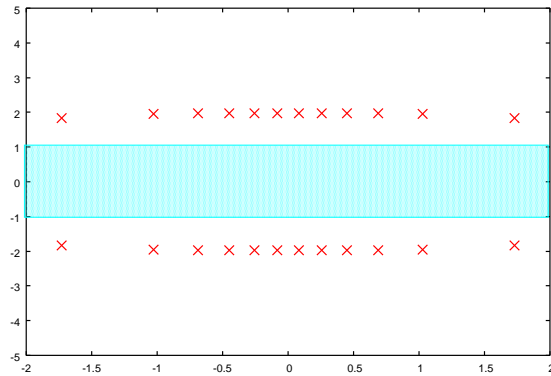
T_1 (left) T_2 (right)



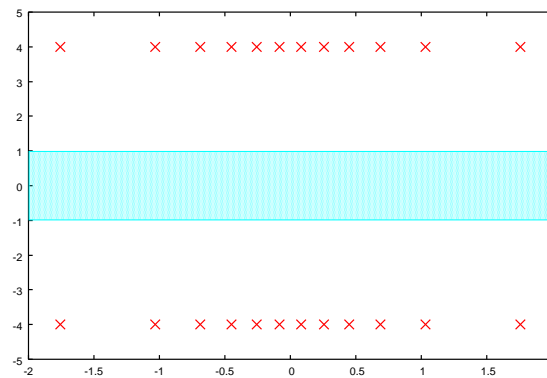
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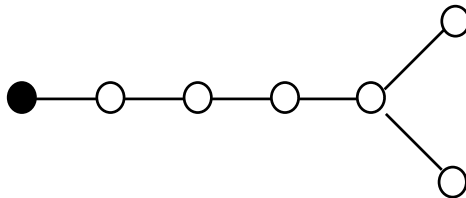
The analyticity and TBA

Analyticity conjecture:

- $Y_j(j = 1, \dots, p - 2, t)$ are analytic and nonzero in the strip $\Im z \in [-1, 1]$ (except for trivial poles)
- $1 + Y_j(j = 1, \dots, p - 2, t)$ are analytic and nonzero in the strip $\Im z \in [-0^+, 0^+]$

1. Truncated Y system
2. Analytic property

\implies TBA with a massive node



The energy of the spin chain

Suppose one solves TBA and obtains Y . Then the functional relation

$$T_1(x+i)T_1(x-i) = f_1(1 + Y_1(x))$$

and the analyticity determines $T(x)$ in the strip $\Im x \in [-1, 1]$.

The energy of the spin chain can be evaluated, as

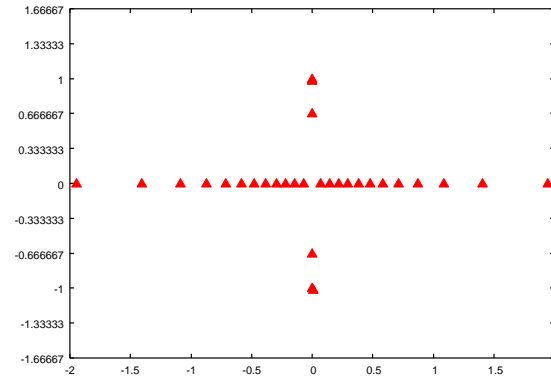
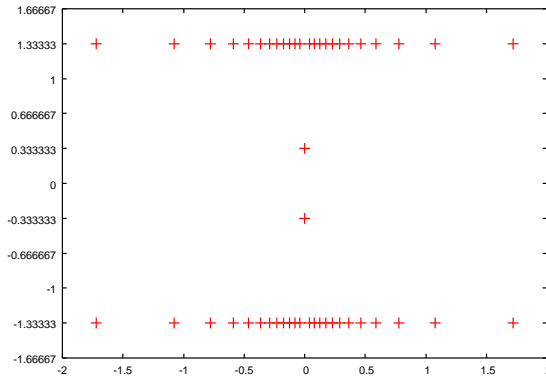
$$E = i \frac{d}{dx} \log T_1(x) \Big|_{x=i}$$

What happens in the attractive regime?

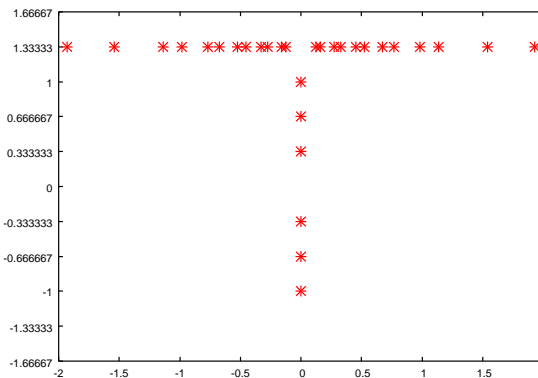
Concentrate on the case $\gamma = 3\pi/4$ with periodicity, $T(x + 8/3i) = T(x)$

Zeros of T in the attractive regime

$$N = 12, p = p' = 0.1, \gamma = \frac{3\pi}{4}$$



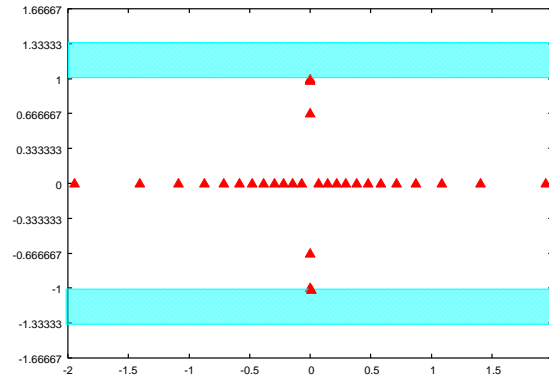
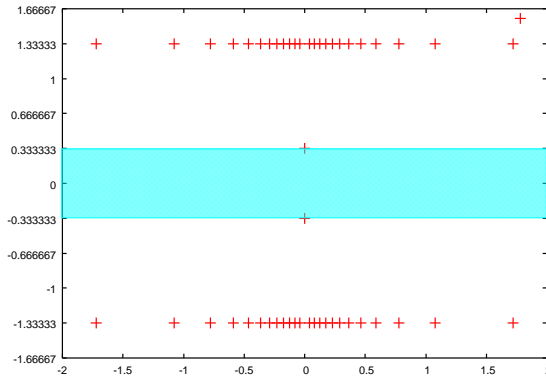
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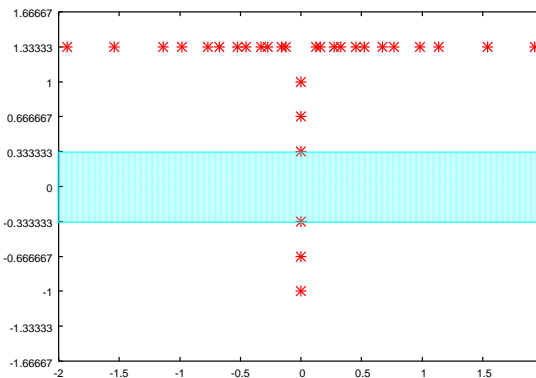
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T_1 (left) T_2 (right)



T_3

Relevant T system

- the original T system

$$T_1(x+i)T_1(x-i) = f_1(x) + T_0(x)T_2(x)$$

$$T_2(x+i)T_2(x-i) = f_2(x) + T_1(x)T_3(x)$$

$$T_3(x+i)T_3(x-i) = f_3(x) + T_2(x)T_4(x)$$

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- the shifted T system Let $T^\vee(x) = T^\vee(x + 4/3i)$

$$T_1^\vee(x+i)T_1^\vee(x-i) = f_1^\vee(x) + T_0^\vee(x)T_2^\vee(x)$$

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$$T_1\left(x + \frac{1}{3}i\right)T_1\left(x - \frac{1}{3}i\right) = f_1^\vee(x) + T_0^\vee(x)T_2^\vee(x)$$

$$T_2\left(x + \frac{1}{3}i\right)T_2\left(x - \frac{1}{3}i\right) = f_2^\vee(x) + T_1^\vee(x)T_3^\vee(x)$$

$$T_3\left(x + \frac{1}{3}i\right)T_3\left(x - \frac{1}{3}i\right) = f_3^\vee(x) + T_2^\vee(x)T_4^\vee(x)$$

Relevant T system

- the original T system

$$T_1^\vee\left(x + \frac{1}{3}i\right)T_1^\vee\left(x - \frac{1}{3}i\right) = f_1(x) + T_0(x)T_2(x)$$

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Relevant Y system

- Define Y functions

$$Y_1(x) = T_0^\vee(x)T_2^\vee(x)/f_1^\vee(x)$$

$$Y_2(x) = T_1(x)T_3(x)/f_2(x)$$

$$Y_t(x) = T_2^\vee(x)/T_0(x)B_-(x - i/3)B_+(x + i/3)$$

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$$Y_1(x + i/3)Y_1(x - i/3) = (1 + Y_2(x))$$

$$Y_2(x + i/3)Y_2(x - i/3) = (1 + Y_1(x))$$

$$\times (1 + \lambda(x)Y_t(x))(1 + \lambda'(x)Y_t(x))$$

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- Y system

$$Y_1(x + i/3)Y_1(x - i/3)/Y_2(x) = (1 + (Y_2(x)))^{-1}$$

$$Y_2(x + i/3)Y_2(x - i/3)/(Y_1(x)Y_t^2(x)) = (1 + (Y_1(x)))^{-1} \\ \times (\lambda(x) + (Y_t(x))^{-1})(\lambda'(x) + (Y_t(x))^{-1})$$

$$Y_t(x + i/3)Y_t(x - i/3)/Y_2(x) = (1 + (Y_2(x)))^{-1}$$

TBA in the attractive regime

- Y system + analyticity assumption
 \implies recovers TBA with the correct mass ratio

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⇒ recovers TBA with the correct mass ratio
- Is the problem solved?

TBA in the attractive regime

- Y system + analyticity assumption
 \implies recovers TBA with the correct mass ratio
- Is the problem solved?
- No. $T_1(x)$ known in $\Im x \in [-1/3, 1/3]$
The evaluation of the energy though

$$E = i \frac{d}{dx} \log T_1(x) \Big|_{x=i}$$

is not yet possible.

A trial?

The analytic continuation of the equation? (i.e., second determination in NLIE.)

- an equation valid in the narrow strip

$$\log T_1(x) = \frac{1}{2\pi} \int K_{\frac{1}{3}}(x - x') \log(1 + Y_1(x')) dx'$$

$$K_{\frac{1}{3}}(x) = \frac{3\pi}{2 \cosh \frac{3\pi}{2}(x)}$$

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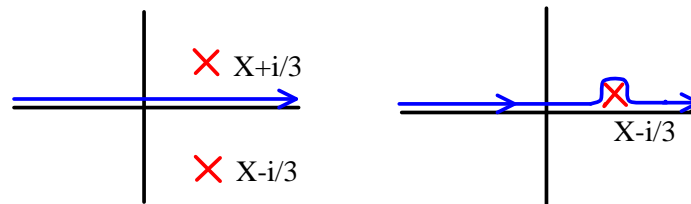
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- analytic continuation ($\Im x = 0$ to $\Im x > 1/3$)



Contour deformation

- Deform the contour to the straightline,

$$\begin{aligned} \log T_1(x + 2/3i) &= \log(1 + Y_1(x + 1/3i)) \\ &\quad - \frac{1}{2\pi} \int_{-\infty}^{\infty} K_{\frac{1}{3}}(x - x') \log(1 + Y_1(x')) dx' \end{aligned}$$

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- Using functional relation

$$\begin{aligned}\log T_1(x + 2/3i) &= \log T_1(x + 2/3i) + \log T_1(x) \\ &\quad - \frac{1}{2\pi} \int K_{\frac{1}{3}}(x - x') \log(1 + Y_1(x')) dx'\end{aligned}$$

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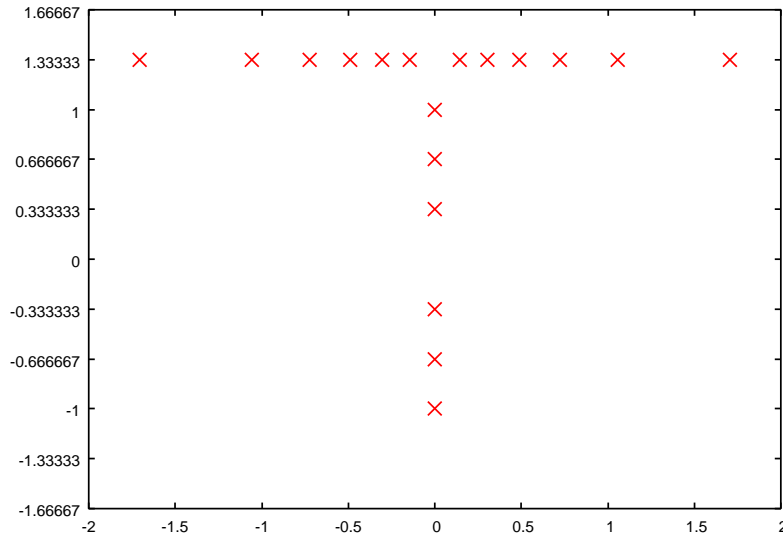
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- Come back to the original form..nonsense!

$$\log T_1(x) = \frac{1}{2\pi} \int K_{\frac{1}{3}}(x - x') \log(1 + Y_1(x')) dx'$$

Zeros of T_3 in the attractive regime

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Factorized T_3

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 $T_3(x) \propto Q(x - 4/3i)Q(x + 4/3i)$.
- for nonzero boundary fields p , let $\bar{Q}(x)$ be Q for a model with $-p$. then

$$T_3(x) \propto Q(x - 4/3i)\bar{Q}(x + 4/3i)$$

•

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$$E \sim i \frac{d}{dx} \log \frac{Q(x - i)}{Q(x + i)} \Big|_{x=0}$$

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$$E \sim i \frac{d}{dx} \log \frac{Q(x - i)}{Q(x + i)} \Big|_{x=0}$$

- Use the factorized T_3

$$E \sim i \frac{d}{dx} \log \frac{T_3(x + i/3)}{T_3(x - i/3)} \Big|_{x=0}$$

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- Interesting to extend this to $\gamma = \pi s / r$ general to see how the factorization plays a role
- THANK YOU